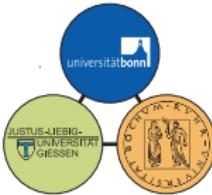


Giessen coupled-channel results for pion and photon induced reactions

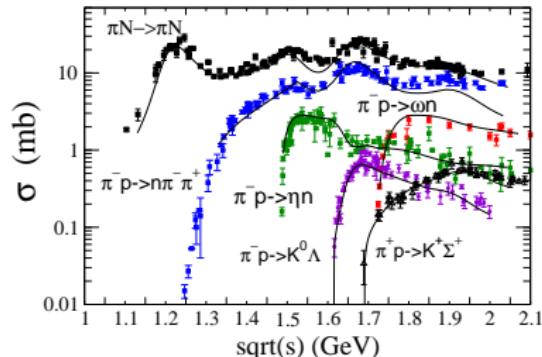
V. Shklyar U. Mosel H. Lenske

Institut für Theoretische Physik
Universität Giessen



Partial wave version of optical theorem

constraints on partial wave cross sections

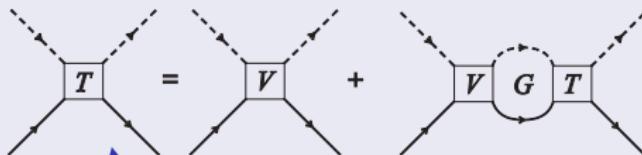


$$\begin{aligned} Im T_{\pi N \rightarrow \pi N}^{JP} = & \frac{k^2}{4\pi} (\sigma_{\pi N \rightarrow \pi N}^{JP} + \sigma_{\pi N \rightarrow 2\pi N}^{JP} + \sigma_{\pi N \rightarrow \eta N}^{JP} \\ & + \sigma_{\pi N \rightarrow \omega N}^{JP} + \sigma_{\pi N \rightarrow K\Lambda}^{JP} + \sigma_{\pi N \rightarrow K\Sigma}^{JP} + \dots) \end{aligned}$$

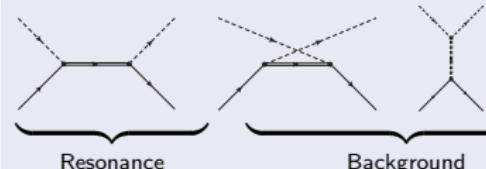
all reaction data are linked
→ need for coupled-channel unitary analysis

Bethe-Salpeter in K -matrix: dynamical model: based on eff. L_{mBB}

T-matrtix



Interaction term V



multidmentional T-matrix

$$T = \begin{pmatrix} T_{\gamma\gamma} & T_{\gamma\pi} & T_{\gamma\eta} & T_{\gamma\omega} & \dots \\ T_{\pi\gamma} & T_{\pi\pi} & T_{\pi\eta} & T_{\pi\omega} & \dots \\ T_{\eta\gamma} & T_{\eta\pi} & T_{\eta\eta} & T_{\eta\omega} & \dots \\ \dots & \dots & \dots & \dots & \dots \end{pmatrix}$$

How many channels?

$\gamma N \rightarrow \gamma N$	$\pi N \rightarrow \pi N$
$\gamma N \rightarrow \pi N$	$\pi N \rightarrow 2\pi N$
$\gamma N \rightarrow \eta N$	$\pi N \rightarrow \eta N$
$\gamma N \rightarrow \omega N$	$\pi N \rightarrow \omega N$
$\gamma N \rightarrow K\Lambda$	$\pi N \rightarrow K\Lambda$
$\gamma N \rightarrow K\Sigma$	$\pi N \rightarrow K\Sigma$

K-matrix approximation:

To solve Bethe-Salpeter equation take the imaginary part of the propagator:

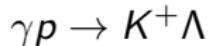
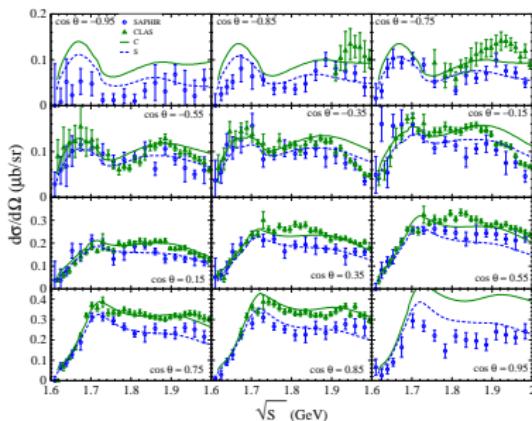
$$\int dq \frac{1}{q^2 - m^2 \pm i\varepsilon} = P \int dq \frac{1}{q^2 - m^2} \mp i\pi \int dq \delta(q^2 - m^2)$$

where all intermediate particles are **on-shell**.

main features

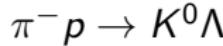
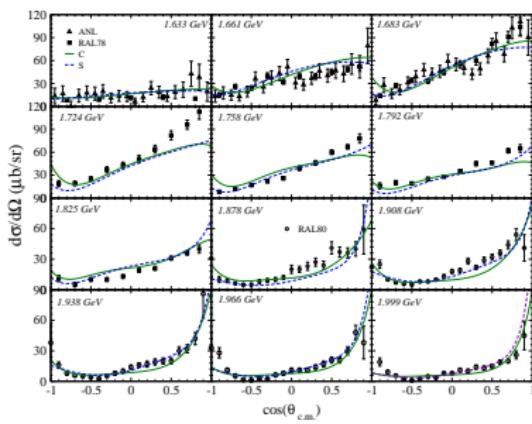
- neglect real part of self energy
- Minkowsky space
- resonance parameters: **coupling constants at interaction
Lagrangians**

$(\gamma, \pi)N \rightarrow K\Lambda$. Giessen model PRC72:015210



Two independent solutions:
C(CLAS) and **S**(SAPHIR)

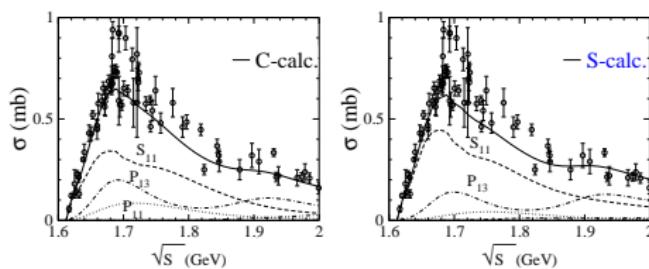
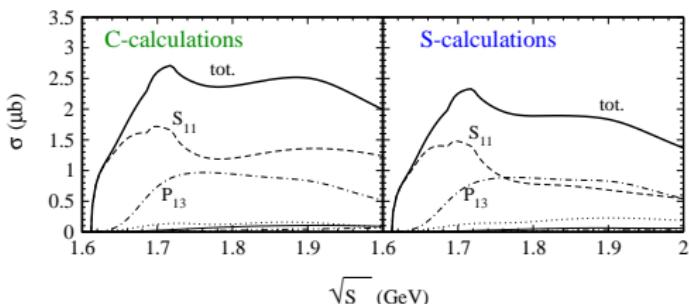
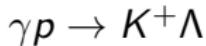
The difference between the C and S-calculations is mostly due to non-resonance contributions.
 (next transp.)



Disagreement between the CLAS and SAPHIR data does not affect the the $\pi^- p \rightarrow K^0\Lambda$ reaction.

$K\Lambda$ -production. Reaction mechanism

Giessen PRC72, 015210 (2005).



Resonance contributions: $S_{11}(1650)$, $P_{13}(1720)$ and $P_{13}(1900)$

$L_{2I,2S}$	$R_{K\Lambda}(C)$	$R_{K\Lambda}(S)$
$S_{11}(1650)$	3.2(+)	4.6(+)
$P_{13}(1720)$	4.6(+)	4.0(+)
$P_{13}(1900)$	2.4(+)	2.3(+)

Table: N^* decay ratios to $K\Lambda$

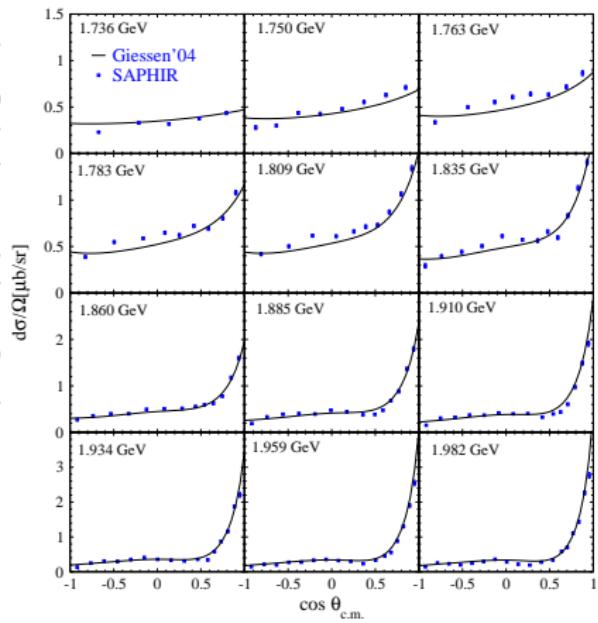
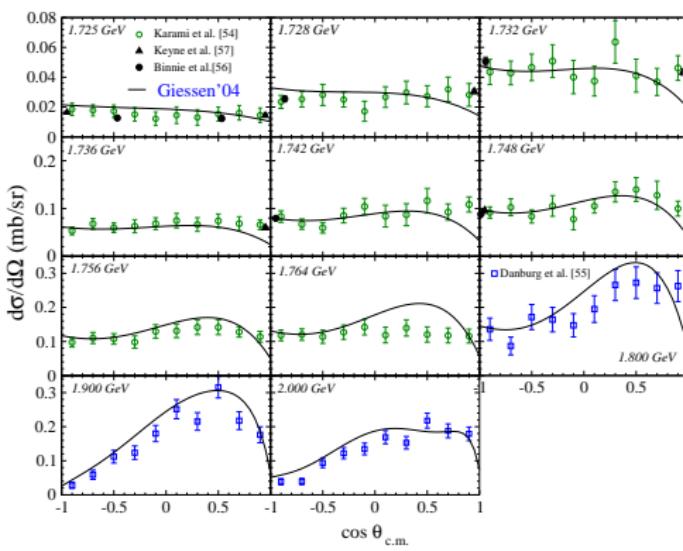
Giessen model. Results for the $(\pi, \gamma)N \rightarrow \omega N$ reactions

Combined analysis \Rightarrow more constraint on resonance properties.

Giessen model, PRC 71:055206, 2005

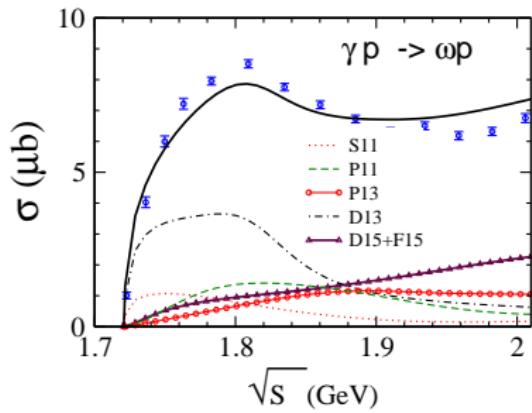
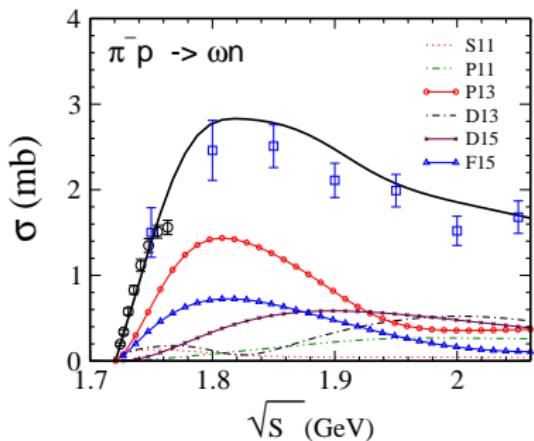
$$\pi N \rightarrow \omega N$$

$$\gamma N \rightarrow \omega N$$



Giessen model. Results for $(\pi, \gamma)N \rightarrow \omega N$

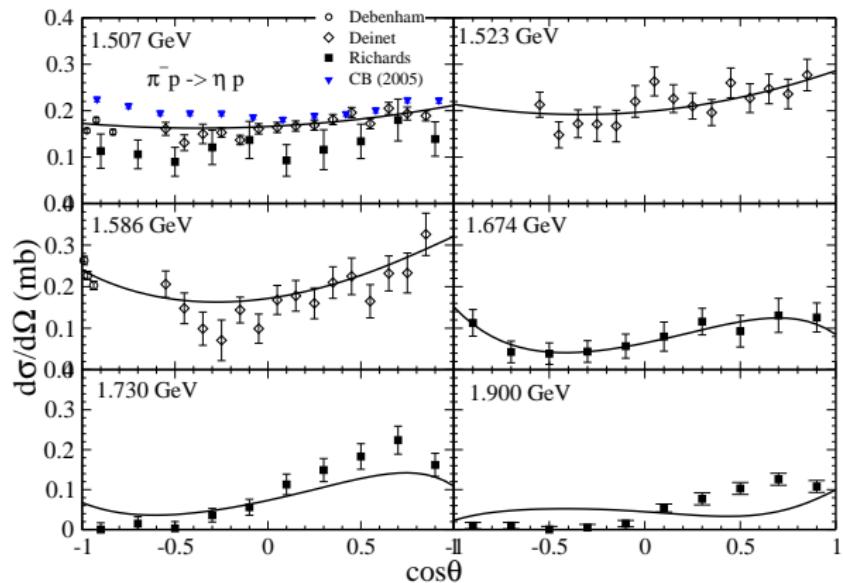
Giessen model, PRC 71:055206, 2005



- P_{13} : interference between resonance and background
- strong $N^*(\frac{5}{2})$ coupling to ωN
- $D_{13}(1520)$ minor contributions

- strong Born and π^0 -exchange contributions
- D_{13} is due to π^0 -exchange

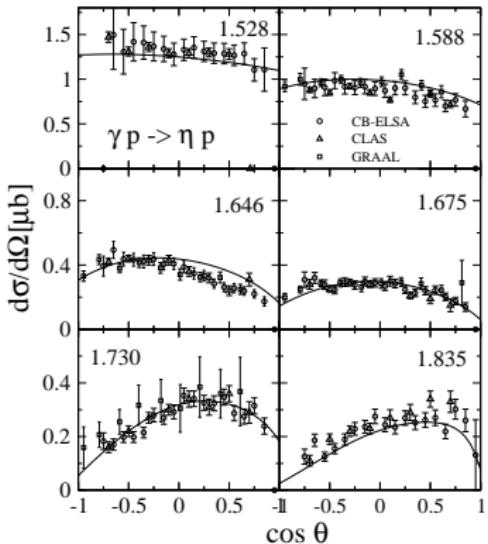
Results for the $\pi^- p \rightarrow \eta p$ production



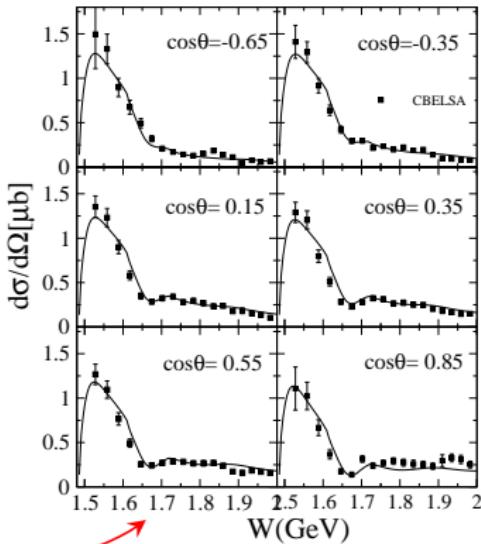
$\pi^- p \rightarrow \eta n$: Solution from the Giessen coupled-channel analysis
V.Shklyar et al, PRC.71. 055206 (2005).

Results for the $\gamma p \rightarrow \eta p$

$\frac{d\sigma}{d\Omega}$ as a function of $\cos(\theta)$



$\frac{d\sigma}{d\Omega}$ as a function of W

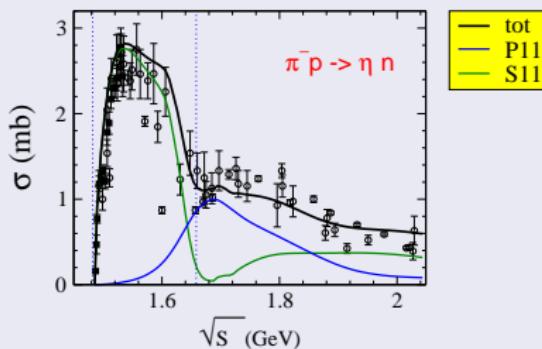
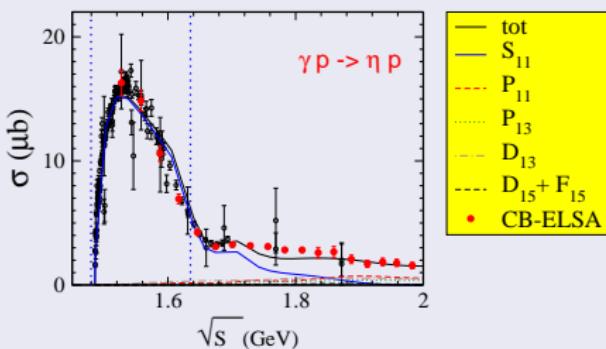


The structure at 1.67 GeV in $\gamma p \rightarrow \eta p$ is due to $S_{11}(1650)$
Shklyar et al PLB650, 172(2007)

no need for any exotic state!

Giessen model

$S_{11}(1535)$ dominates
both $\gamma p \rightarrow \eta p$ and $\pi^- p \rightarrow \eta n$ reactions

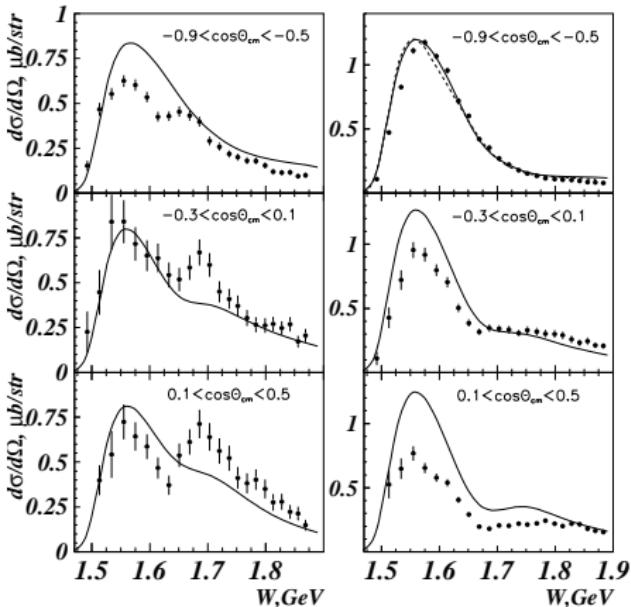


- strong $S_{11}(1535)$ excitation
- kink structure at 1.72 GeV is due to the ωN threshold
- seems no room for other contributions
- destructive effect from $S_{11}(1650)$
- above 1.6 GeV - $P_{11}(1710)$ - consistent with πN inelasticity

$$\gamma n^* \rightarrow \eta n$$

V. Kuznetsov, et al. PLB 647 (2007) 23 for GRAAL collaboration

$$\gamma n^* \rightarrow \eta n \quad \gamma p^* \rightarrow \eta p$$



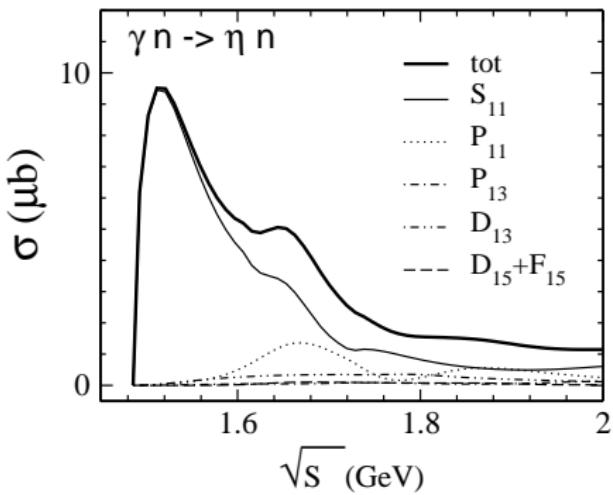
- quasi-free neutron: resonance-like structure at 1.67 GeV
- confirmed by B.Krusche, I. Jaegle at MAMI, CB-ELSA

Possible explanations

- Polyakov, Strakovsky, Arndt, Workman; Polyakov Kuznetsov: pentaquark partner
- Shklyar, Mosel, Lenske: well known $S_{11}(1650)$, $P_{11}(1710)$
- M. Doering: cusp in $K\Sigma$

Results for the $\gamma n \rightarrow \eta n$

Giessen Model PLB650, 172(2007): total $\gamma n \rightarrow \eta n$ cross section

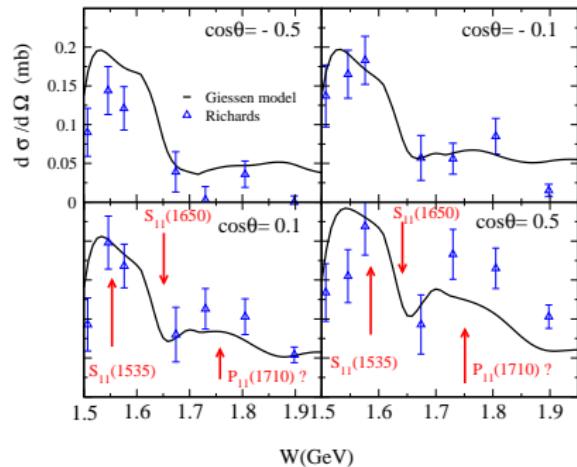


$$A_{1/2}^n(1650) = -9 \times 10^{-3} \text{ GeV}^{-\frac{1}{2}}$$

$$A_{1/2}^n(1710) = 24 \times 10^{-3} \text{ GeV}^{-\frac{1}{2}}$$

$$\pi^- p \rightarrow \eta n$$

Giessen Model: Shklyar, Mosel, Lenske PLB650, 172(2007)
 vs. data Richards et al PR 1, 10 (1970)

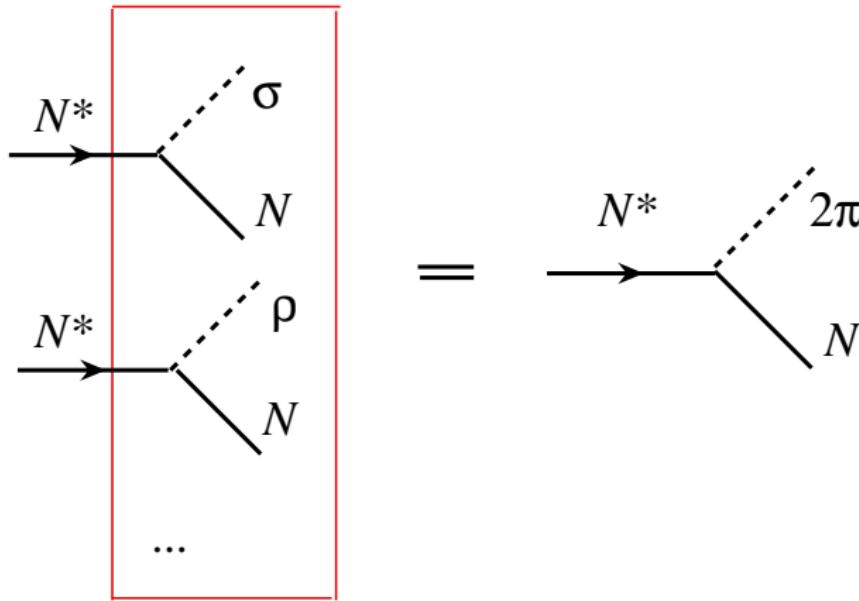


- Richards data show an excess structure at 1.7 GeV
- hard to make conclusion: the data is of poor quality
- Giessen calculations: destructive $S_{11}(1535)$ and $S_{11}(1650)$ interference; $P_{11}(1710)$ excitation.

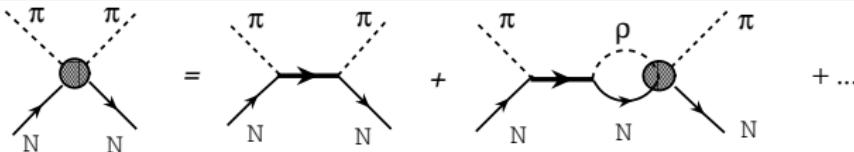
Next step: improve description of the $2\pi N$ channel

so far: N^* decay into 'generic' 2π channel

- take $2\pi N$ inelastic flux into account
- $N^* \rightarrow 2\pi N$ couplings constrained by $\sigma_{\pi N \rightarrow 2\pi N}^{JI}$



New multichannel problem



$$T_{\pi\pi}^{\text{JI}} = K_{\pi\pi}^{\text{JI}} + i K_{\pi\pi}^{\text{JI}} T_{\pi\pi}^{\text{JI}}$$

$$+ i \int_{4m_\pi^2}^{(\sqrt{s}-m_N)^2} d\mu_\rho'^2 K_{\pi\rho}^{\text{JI}}(\mu_\rho'^2) A_\rho(\mu_\rho'^2) T_{\rho\pi}^{\text{JI}}(\mu_\rho'^2)$$

summation instead of integration

$$T_{\pi\pi}^{\text{JI}} = K_{\pi\pi}^{\text{JI}} + i K_{\pi\pi}^{\text{JI}} T_{\pi\pi}^{\text{JI}}$$

$$+ i \sum_{m_{\rho_i}} 2m_{\rho_i} \Delta m_{\rho_i} K_{\pi\rho_i}^{\text{JI}}(m_{\rho_i}^2) A_{\rho_i}(m_{\rho_i}^2) T_{\rho_i\pi}^{\text{JI}}(m_{\rho_i}^2)$$

$N(1520)$ D_{13} state

Manley et al: PRD(1984)

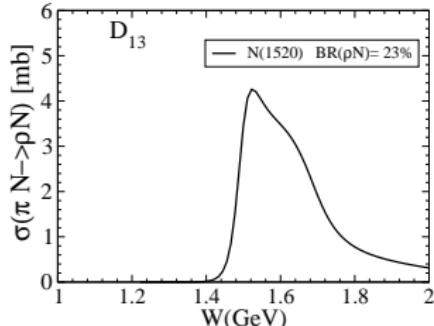
$M_R = 1520\text{MeV}$

$\Gamma_{\text{tot}} = 120\text{MeV}$

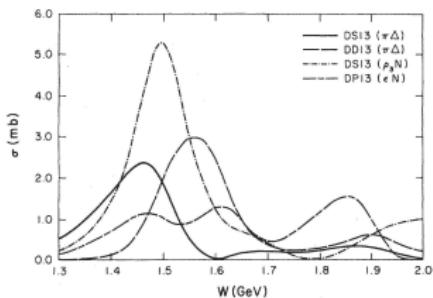
strong $N(1520) \rightarrow 2\pi N$

$\text{Br}(\rho N) \approx 20\%$

Giessen Model (CC): $\pi N \rightarrow \rho N$



Manley analysis:



- distribution:
Giessen: **non-symmetric**
Manley : **symmetric**
- Gi Model: no contributions below 1.4 GeV
- Manley: no ρ -spectral function: **should be revised**

$$\pi N \rightarrow 2\pi N$$

Summary of the $\pi N \rightarrow 2\pi N$ reactions

- strong contributions to the πN inelasticity
- important for understanding for ρ -meson dynamics and resonance couplings
- could solve many puzzles in non-strange baryon spectroscopy: origin and properties of the $P_{11}(1440)$, $P_{11}(1710)$, $D_{13}(1520)$ etc.

Theory

- analysis of Manley et. al. should be revised!

Experiment

- need for new measurements $\pi N \rightarrow 2\pi N$ in region 1.2...2. GeV → challenge for HADES collaboration

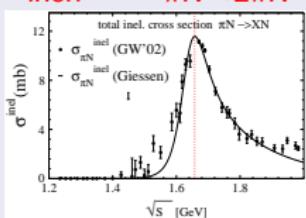
pion beams at HADES contact piotr.salabura@uj.edu.pl

Why $D_{15}(1675)$ with $\Gamma_{\eta N} = 17\%$ is a bad

Optical theorem for $\pi N \rightarrow \pi N$ scattering

$$(J+\frac{1}{2}) Im T_{\pi N \rightarrow \pi N}^{\frac{5}{2}+\frac{1}{2}} = \frac{k^2}{4\pi} (\sigma_{\pi N \rightarrow \pi N}^{\frac{5}{2}+\frac{1}{2}} + \sigma_{\pi N \rightarrow 2\pi N}^{\frac{5}{2}+\frac{1}{2}} + \sigma_{\pi N \rightarrow \eta N}^{\frac{5}{2}+\frac{1}{2}})$$

$$\sigma_{\text{inel.}}^{\frac{5}{2}+\frac{1}{2}} = \sigma_{\pi N \rightarrow 2\pi N}^{\frac{5}{2}+\frac{1}{2}} + \sigma_{\pi N \rightarrow \eta N}^{\frac{5}{2}+\frac{1}{2}} \approx 12 \text{mb}$$



Assuming dominant $D_{15}(1675)$

$$T_{\pi N \rightarrow 2\pi N} \sim \frac{\Gamma_{\pi N}^{N(1675)} \Gamma_{2\pi N}^{N(1675)}}{s - m_{1675}^2 - i\Gamma_{\text{tot}}/2}$$

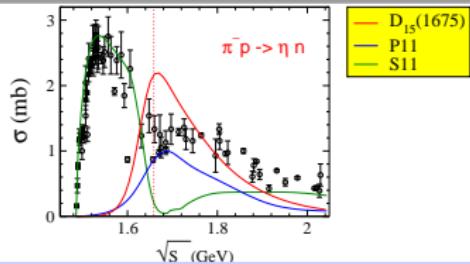
$$T_{\pi N \rightarrow \eta N} \sim \frac{\Gamma_{\pi N}^{N(1675)} \Gamma_{\eta N}^{N(1675)}}{s - m_{1675}^2 - i\Gamma_{\text{tot}}/2}$$

$$\sigma_{\pi N \rightarrow \eta N} / \sigma_{\pi N \rightarrow 2\pi N} = (\Gamma_{\eta N}^{1675} / \Gamma_{2\pi N}^{1675})^2$$

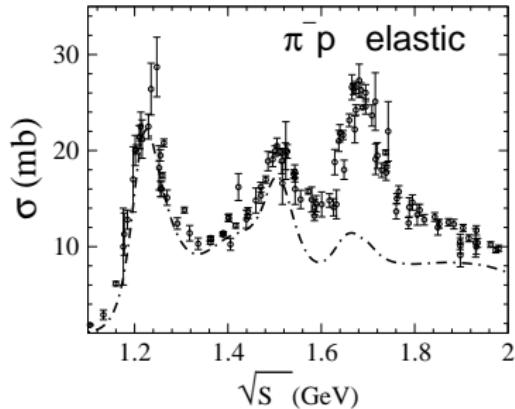
η -MAID L. Tiator (hep-ex/0601002):

$$\Gamma_{\eta N} \approx 17\%, \Gamma_{2\pi N} \approx 40\%$$

$$\sigma_{\pi N \rightarrow \eta N}^{\frac{5}{2}}|_{1675 \text{ MeV}} = 12 \left(\frac{17}{40}\right)^2 \approx 2.2 \text{ mb}$$

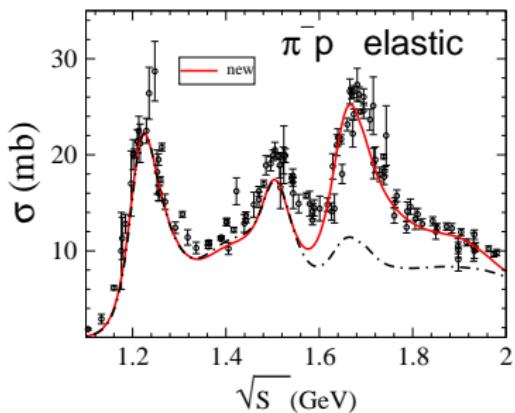


Giessen model



Previous analysis:
Penner and Mosel RRC66,
055211 (2002)
no spin- $\frac{5}{2}$ resonances !

Giessen model



New results:

V. Shklyar et al .PRC71,
055206 (2005)

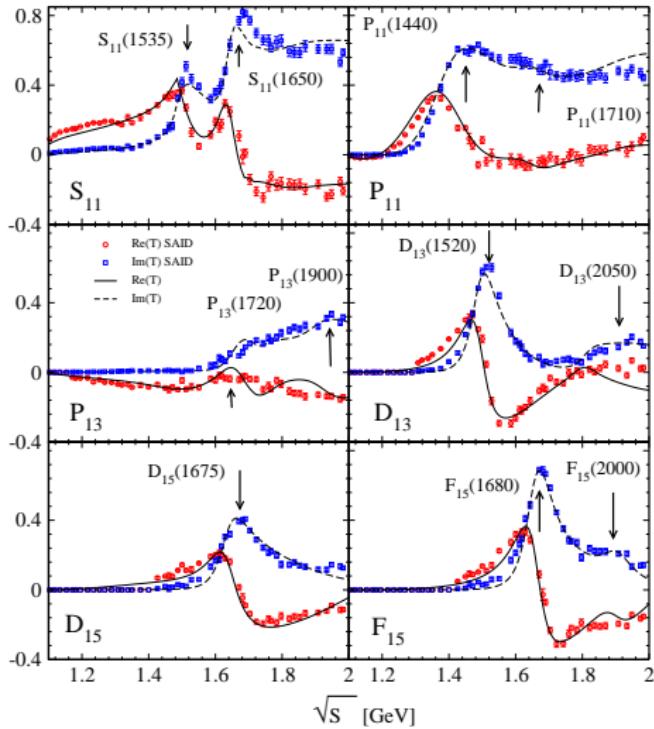
with spin- $\frac{5}{2}$ resonances !

But! It is so important for the
 ωN production ?

Optical theorem:

$$Im T_{\pi N \rightarrow \pi N} \sim \sigma_{\pi N \rightarrow \omega N} + \dots$$

Results for pion-induced reactions



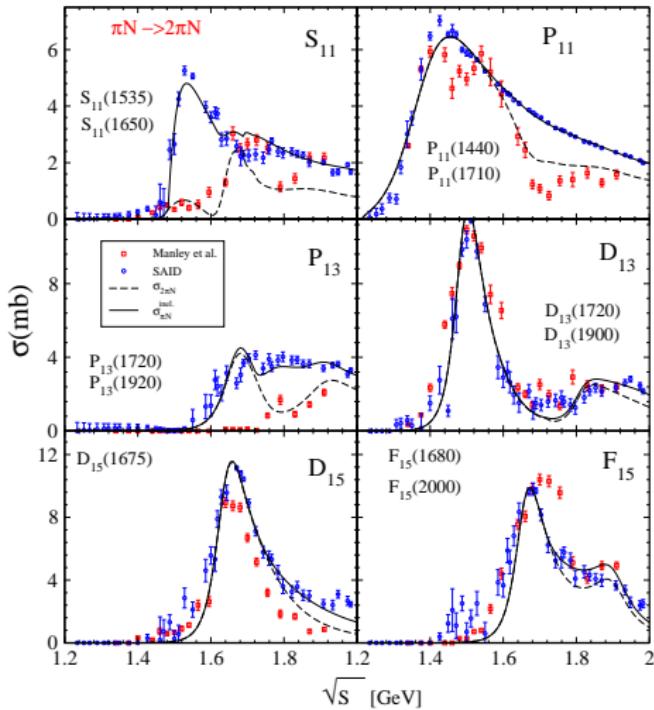
πN elastic amplitudes:

- reach spectrum
- not only πN decay

$I = \frac{1}{2}$ resonances important:

- $S_{11}(1535), S_{11}(1650)$
- $P_{11}(1440), P_{11}(1710)$
- $P_{13}(1720), P_{13}(1900)$
- $D_{13}(1520), D_{13}(2050)$
- $D_{15}(1675)$
- $F_{15}(1680), F_{15}(2000)$

πN inelasticity and inelastic channels



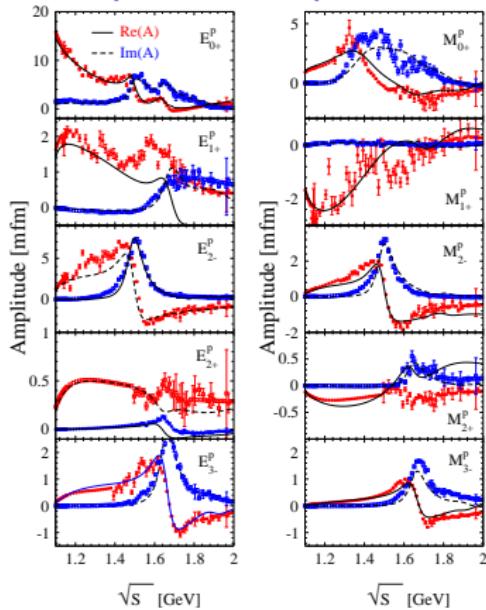
Optical theorem :

$$\begin{aligned}
 & \left[\frac{4\pi}{k_{cm}^2} \text{Im} T_{\pi N}^{JI} - \sigma_{\pi N \rightarrow \pi N}^{JI} \right] \\
 &= \sigma_{\pi N \rightarrow 2\pi N}^{JI} + \sigma_{\pi N \rightarrow \eta N}^{JI} \\
 &+ \sigma_{\pi N \rightarrow \omega N}^{JI} + \sigma_{\pi N \rightarrow K\Lambda}^{JI} + \sigma_{\pi N \rightarrow K\Sigma}^{JI}
 \end{aligned}$$

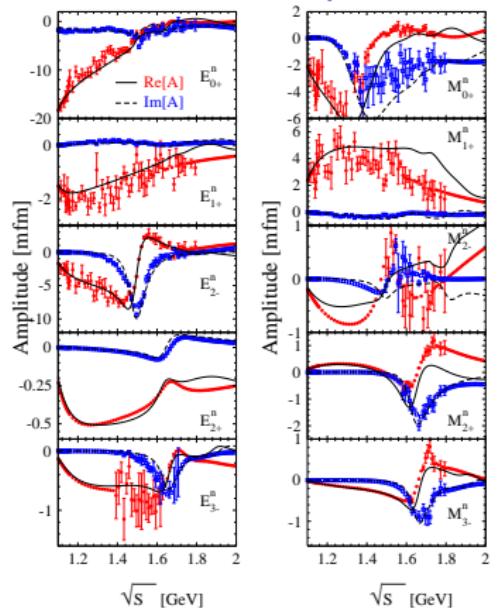
— πN inelasticity
— $2\pi N$ partial wave cross sections

Giessen model. Pion photoproduction

proton multipoles



neutron multipoles



Combined analysis of $(\pi, \gamma)N \rightarrow (\pi, \gamma)N$ gives a strong constraint on extracted resonance parameters